

# Inelastic RC Systematic Error

Ryan Zielinski

4/6/16

# Overview

- Want to do the same thing I did for the elastic tail and check the error with the approximations and calculations done in inelastic RCs
- Some errors carry over from the elastic tail so I don't have to redo them:
  - $F_{\text{bar}}$  (higher order virtual photon corrections)
  - Energy-peaking approximation used in evaluating external corrections
  - Born approximation
- New potential sources of error
  - Angle-peaking approximation in internal bremsstrahlung
  - Choice of straggling function for external bremsstrahlung
    - Did check this for the elastic tail but should do it again
  - Soft photon correction factor
  - Interpolation/extrapolation error
    - Including using a model as source for lowest energy setting to RADCOR input

# Internal Bremsstrahlung

- For the elastic tail, the full expression for the internal bremsstrahlung is

$$\begin{aligned}
 \sigma_{\text{exact}} = & \left( \frac{d^2\sigma}{d\Omega dE_p} \right)_{\text{ex}} = \frac{\alpha^3}{2\pi} \left( \frac{E_p}{E_s} \right) \int_{-1}^1 \frac{2M_T\omega d(\cos\theta_k)}{q^4(u_0 - |\vec{u}|\cos\theta_k)} \\
 & \times \left( \tilde{W}_2(q^2) \left\{ \frac{-am^2}{x^3} \left[ 2E_s(E_p + \omega) + \frac{q^2}{2} \right] - \frac{a'm^2}{y^3} \left[ 2E_p(E_s + \omega) + \frac{q^2}{2} \right] \right. \right. \\
 & - 2 + 2\nu(x^{-1} - y^{-1}) \{ m^2(s \cdot p - \omega^2) + (s \cdot p)[2E_sE_p - (s \cdot p) + \omega(E_s - E_p)] \} \\
 & + x^{-1} \left[ 2(E_sE_p + E_s\omega + E_p^2) + \frac{q^2}{2} - (s \cdot p) - m^2 \right] \\
 & \left. \left. - y^{-1} \left[ 2(E_pE_s + E_p\omega + E_s^2) + \frac{q^2}{2} - (s \cdot p) - m^2 \right] \right\} \right. \\
 & + \tilde{W}_1(q^2) \left[ \left( \frac{a}{x^3} + \frac{a'}{y^3} \right) m^2(2m^2 + q^2) + 4 + 4\nu(x^{-1} - y^{-1})(s \cdot p)(s \cdot p - 2m^2) \right. \\
 & \left. \left. + (x^{-1} - y^{-1})(2s \cdot p + 2m^2 - q^2) \right] \right),
 \end{aligned}$$

# Angle-Peaking Approximation

- Full integral is calculated at all possible photon emission angles
  - But in reality, the majority of photon's are emitted in the same direction as the incoming and outgoing electrons
- Can then approximate full integral as
  - Uses an 'equivalent radiator' == internal bremsstrahlung approximated as an external bremsstrahlung correction

$$\sigma_{\text{int}}^{\text{pk}}(E_s, E_p) = \frac{bt_r \phi(v_p)}{\omega_p} \tilde{F}(q_p^2) \sigma_{\text{el}}(E_s) + \frac{M_T + 2(E_s - \omega_s) \sin^2 \frac{\theta}{2}}{M_T - 2E_p \sin^2 \frac{\theta}{2}} \frac{bt_r \phi(v_s)}{\omega_s} \tilde{F}(q_s^2) \sigma_{\text{el}}(E_s - \omega_s),$$

$$\omega_s = E_s - \frac{E_p}{1 - 2 \frac{E_p}{M_T} \sin^2 \frac{\theta}{2}},$$

$$\omega_p = \frac{E_s}{1 + 2 \frac{E_s}{M_T} \sin^2 \frac{\theta}{2}} - E_p,$$

$$bt_r = \frac{\alpha}{\pi} \left[ \ln \left( \frac{q^2}{m_e^2} \right) - 1 \right]$$

Energy of incoming (outgoing) photon

Follows Eqs (B.6) – (B.11) of Mo and Tsai, Rev. Mod. Phys 41 205, (1969)

# Inelastic Internal Bremsstrahlung

- Inelastic isn't limited to  $W = M_T$  so the inelastic internal contribution is the elastic internal tail integrated over all possible  $M_f$

$$\frac{d\sigma}{d\Omega dE}(\omega > \Delta) = \frac{\alpha^3}{2\pi} \frac{E_p}{E_s M_T} \int_{-1}^1 d(\cos\theta_k) \int_{\Delta}^{\omega_{\max}(\cos\theta_k)} \frac{\omega d\omega}{q^4} B_{\mu\nu} T^{\mu\nu}$$

Contains the inelastic structure functions

- $\Delta$  parameter avoids a divergence at  $\omega = 0$ , so the full internally inelastically radiated spectrum is

$$\frac{d\sigma_r}{d\Omega dE}(E_s, E_p) = e^{-\delta_r(\Delta)} \tilde{F}(q^2) \frac{d\sigma}{d\Omega dE} + \frac{d\sigma_r}{d\Omega dE}(\omega > \Delta)$$

- This is a slightly tweaked version of (B.6) of MT, but it keeps choice of  $f_{\text{bar}}$  consistent across calculations (See B. Badalek *et al.* arxiv:hep-ph/940328v1 (1994).)

# Inelastic Peaking Approximation

- The equivalent correction in the angle-peaking approximation is
  - Dropping the soft-photon terms. This is SLAC-PUB-848 (1971) (C.23)

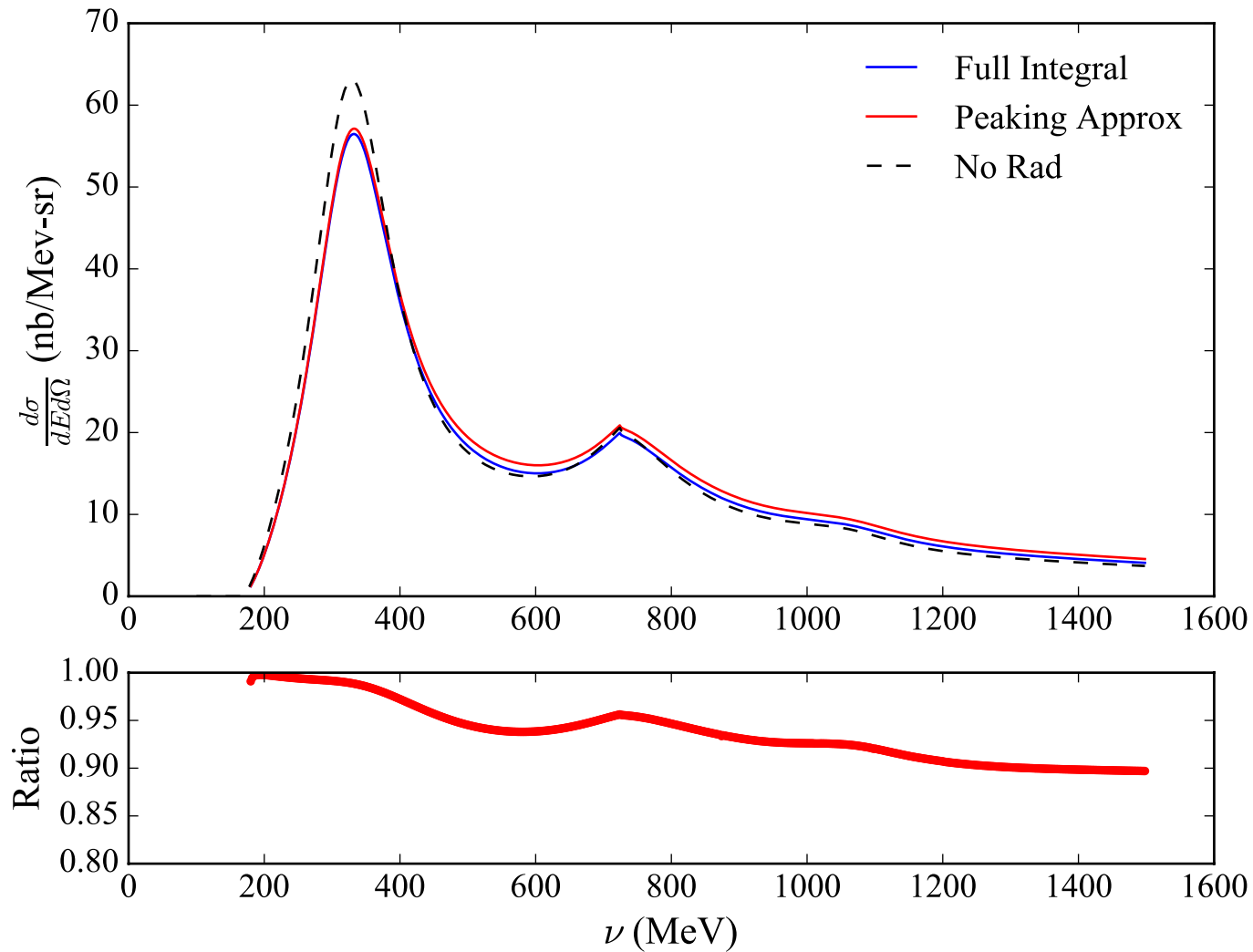
$$\begin{aligned} \frac{d\sigma_r}{d\Omega dE}(E_s, E_p) = & \tilde{F}(q^2) \left[ \left( \frac{R\Delta}{E_s} \right) \left( \frac{\Delta}{E_p} \right)^{bt_r} \frac{d\sigma}{d\Omega dE}(E_s, E_p) \right. \\ & + \int_{E_p+\Delta}^{E_{p\max}} dE'_p \frac{d\sigma}{d\Omega dE}(E_s, E'_p) \frac{bt_r}{2(E'_p - E_p)} \phi\left(\frac{E'_p - E_p}{E'_p}\right) \\ & \left. + \int_{E_{s\min}}^{E_s - R\Delta} dE'_s \frac{d\sigma}{d\Omega dE}(E'_s, E_p) \frac{bt_r}{2(E_s - E'_s)} \phi\left(\frac{E_s - E'_s}{E_s}\right) \right] \end{aligned}$$

- Angle peaking approximation is used because:
  - Significantly faster (<1 min to run compared to multiple hours for full integral)
  - In general the inelastic structure functions aren't known but the total cross section is, so this can be used for unfolding!

# Comparing the Two

Using the Bosted model

p: 2135 Mev/ 6 degrees



# Going Forward

- Difference between the two calculations is 5-10% for the proton
  - Haven't checked yet to see how much internal radiation contributes to total radiated cross section
- Need to make the switch to nitrogen for the saGDH data
- Machinery is also in place to check model input systematic
- Already looked at the systematic for interpolation
  - Previously showed a model radiated and then unfolded.
- Move on to the other systematic checks I mentioned